28 November 2011 – Tallinn

THE NUMERICAL COMPUTATION OF VIOLENT LIQUID MOTION



Frédéric DIAS

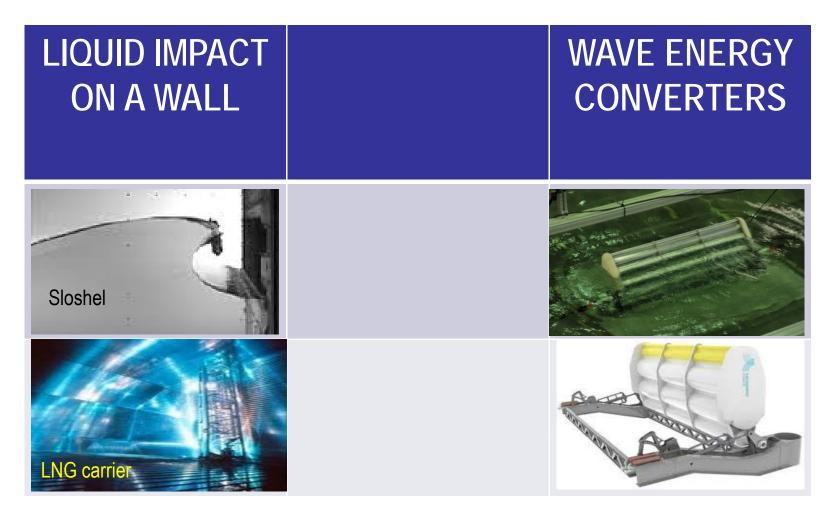
School of Mathematical Sciences University College Dublin on leave from Ecole Normale Supérieure de Cachan







TOPICS COVERED IN TODAY'S TALK



COLLABORATORS

- ✤Jean-Philippe Braeunig (INRIA & CEA)
- Laurent Brosset (GTT)
- Paul Christodoulides (Cyprus University of Technology)
- ✤Ken Doherty (Aquamarine Power Ltd.)
- Christophe Fochesato (CEA)
- ✤Jean-Michel Ghidaglia (Ecole Normale Supérieure de Cachan)
- Emmanuel Reynaud (University College Dublin)

Postdocs

- Emiliano Renzi (University College Dublin)
- Roxana Tiron (University College Dublin)

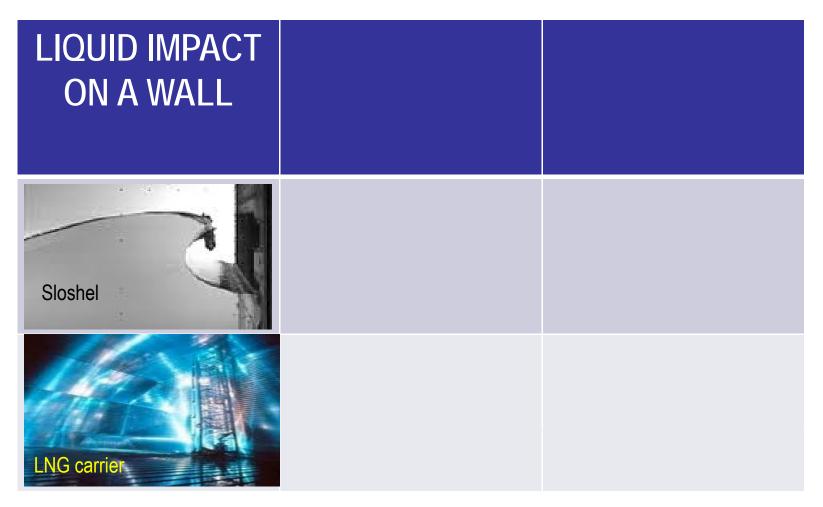
PhD students

- Laura O'Brien (University College Dublin)
- Sarah Gallagher (University College Dublin)

Research funded by

□SFI (Science Foundation Ireland) □GTT (Gaz Technigaz & Transport)

PART 1



WAVE IMPACT AND PRESSURE LOADS

•Local phenomena involved during wave impacts are very sensitive to input conditions

•The density of bubbles, the local shape of the free surface, the local flow make the impact pressure change dramatically even for the same experimental conditions

•How does one extrapolate wave impact from model (small scale) to prototype (full scale)?

•In recent years, we have addressed the scaling issue by studying the various local phenomena present in wave impact one after the other in order

to better understand the physics behind

to improve the experimental modelling



VIDEO 1 – Experiments in Marseille (courtesy of O. Kimmoun)

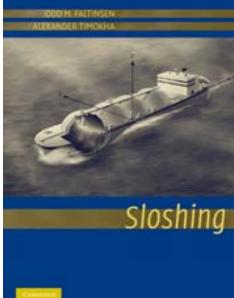
SLOSHING IN TANKS OF LIQUEFIED NATURAL GAS (LNG) CARRIERS





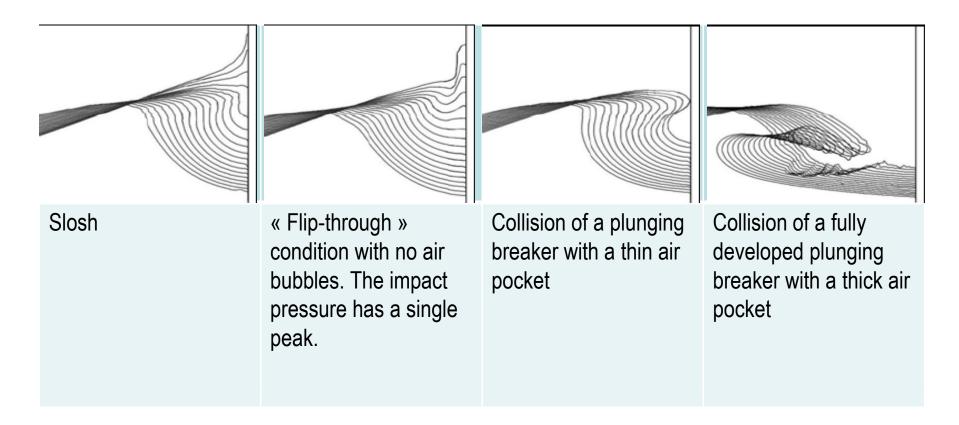
Heavy weather and partially loaded LNG tank resulted sloshing damage in a LNG tank – 6 months and Millions \$\$





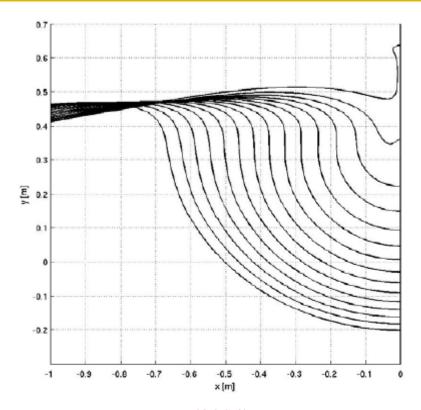


VARIOUS SCENARIOS OF WAVE IMPACT

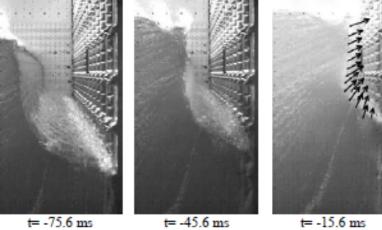


Source : Peregrine D.H. (2003), Water wave impact on walls, Annual Review of Fluid Mechanics

THE FLIP-THROUGH PHENOMENON



Run-up of wave trough
Forward motion of almost vertical wave front
Jet flow generated at the wall 8 m/s



55 bars

Bredmose et al. (2004), Water wave impact on walls and the role of air, *Proceedings ICCE 2004*

Brosset et al. (2011), A Mark III Panel Subjected to a Flip-through Wave Impact: Results from the Sloshel Project, *Proceedings ISOPE 2011*

PHENOMENOLOGY OF A LIQUID IMPACT

Global behaviour

Global flow governed by Froude number

Local behaviour

Escape of the gas between the liquid and the wall: momentum transfer between liquid and gas

Compression of the partially entrapped gas during the last stage of the impact

➢ Rapid change of momentum of the liquid diverted by the obstacle

➢ Possible creation of shock waves: pressure wave within the liquid and strain wave within the wall

> Hydro-elasticity effects during the fluid-structure interaction

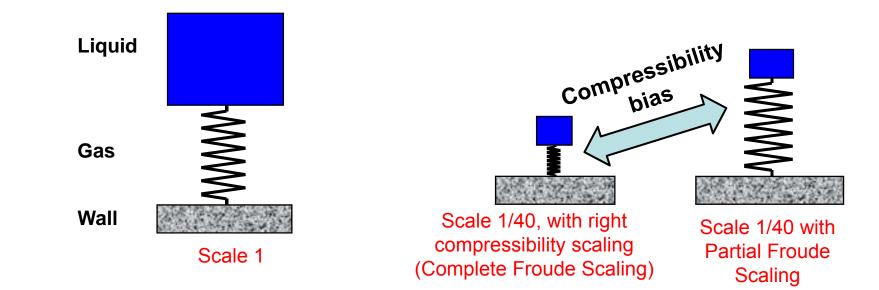
Braeunig et al. (2009), Phenomenological Study of Liquid Impacts through 2D Compressible Two-fluid Numerical Simulations, *Proceedings ISOPE 2009*



EXAMPLE OF BIAS INTRODUCED AT SMALL SCALE

Locally, the impact process is not similar for similar inflow conditions

- Gas compressibility bias: the equations of state should be scaled
- > Liquid compressibility bias: speed of sound should be scaled



Biases are different for different impacts: no unique scaling law!

INCOMPRESSIBLE EULER EQUATIONS WITH AN INTERFACE IN PHYSICAL VARIABLES

Fluid equations:

$$u_{t} + uu_{x} + vu_{y} + wu_{z} + \frac{1}{\rho}\frac{\partial p}{\partial x} = 0$$

$$u_{t} + uv_{x} + vv_{y} + wv_{z} + \frac{1}{\rho}\frac{\partial p}{\partial y} = 0$$

$$w_{t} + uw_{x} + vw_{y} + ww_{z} + \frac{1}{\rho}\frac{\partial p}{\partial z} = -g$$

$$u_{x} + v_{y} + w_{z} = 0$$

$$g: \text{ acceleration due to gravity}$$

Boundary conditions:

h(x, y, t): elevation of the interface

(x, y, z): spatial coordinates

$$\begin{array}{rcl} \mathbf{u} \cdot \mathbf{n} &=& 0 \\ h_t + uh_x + vh_y &=& w \\ p &=& 0 \end{array} \end{array} \quad \begin{array}{c} \text{kinematic and dynamic} \\ \text{conditions on the interface} \end{array}$$

INVARIANCE OF INCOMPRESSIBLE EULER EQUATIONS WITH AN INTERFACE

Froude scaling
$$\frac{u_{fs}^2}{gD_{fs}} = \frac{u_{ms}^2}{gD_{ms}}$$

fs stands for full scale, ms for model scale

 $D_{fs} = \lambda D_{ms}$

$$\mathbf{u}_{fs} \left(\lambda \mathbf{x}_{ms}, \sqrt{\lambda} t_{ms} \right) = \sqrt{\lambda} \mathbf{u}_{ms} \left(\mathbf{x}_{ms}, t_{ms} \right) \qquad t_{fs} = \sqrt{\lambda} t_{ms}$$
$$\dot{\mathbf{u}}_{fs} \left(\lambda \mathbf{x}_{ms}, \sqrt{\lambda} t_{ms} \right) = \dot{\mathbf{u}}_{ms} \left(\mathbf{x}_{ms}, t_{ms} \right) \qquad p_{fs} = \lambda p_{ms}$$

TWO-FLUID COMPRESSIBLE EULER EQUATIONS WITH AN INTERFACE IN PHYSICAL VARIABLES

Fluid equations

$$\frac{\partial \rho_k}{\partial t} + \nabla \cdot (\rho_k \vec{u}_k) = 0 \,,$$

$$\rho_k \left(\frac{\partial \vec{u}_k}{\partial t} + (\vec{u}_k \cdot \nabla) \, \vec{u}_k \right) + \nabla p_k = \rho_k (\vec{g} - \vec{\gamma}_e - 2\vec{\Omega} \times \vec{u}_k) \,, \qquad k = 1, 2$$

$$\rho_k \left(\frac{\partial e_k}{\partial t} + (\vec{u}_k \cdot \nabla) \, e_k \right) + p_k \nabla \cdot \vec{u}_k = 0 \,,$$

Boundary conditions

$$\vec{u}_k \cdot \vec{n} = 0$$

$$\frac{\partial f}{\partial t} = -\vec{u}_1\cdot\nabla f = -\vec{u}_2\cdot\nabla f\,,\quad \text{on}\quad f(x_1,x_2,x_3,t) = 0\,,$$

$$p_1 = p_2$$
, $T_1 = T_2$, on $f(x_1, x_2, x_3, t) = 0$.

Equation of state

$$\rho_k = \mathcal{R}_k(p_k, e_k) \,.$$

INVARIANCE OF TWO-FLUID COMPRESSIBLE EULER EQUATIONS WITH AN INTERFACE

Froude scaling
$$\frac{u_{fs}^2}{gD_{fs}} = \frac{u_{ms}^2}{gD_{ms}}$$
 (including the speeds of sound)
 $\mathbf{u}_{fs} \left(\lambda \mathbf{x}_{ms}, \sqrt{\lambda} t_{ms} \right) = \sqrt{\lambda} \mathbf{u}_{ms} \left(\mathbf{x}_{ms}, t_{ms} \right)$
 $\dot{\mathbf{u}}_{fs} \left(\lambda \mathbf{x}_{ms}, \sqrt{\lambda} t_{ms} \right) = \dot{\mathbf{u}}_{ms} \left(\mathbf{x}_{ms}, t_{ms} \right)$

Differences with the one-fluid incompressible case

$$\mu = \rho_{\rm ms} / \rho_{\rm fs} \qquad \qquad \mathcal{R}_k^{\lambda,\mu}(p,e) = \mu \mathcal{R}_k \left(\frac{\lambda}{\mu} p, \lambda e\right) \qquad \qquad \frac{\rho_g^{\rm fs}}{\rho_\ell^{\rm fs}} = \frac{\rho_g^{\rm ms}}{\rho_\ell^{\rm ms}}$$

 $p^{\text{fs}} = \lambda(\rho_{\text{lig}}^{\text{fs}}/\rho_{\text{lig}}^{\text{ms}}) p^{\text{ms}}$

Scale the equations of state

Keep the same density ratio

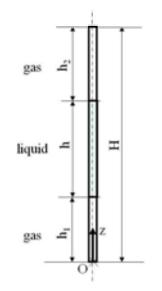
THE ROLE OF NUMERICAL STUDIES

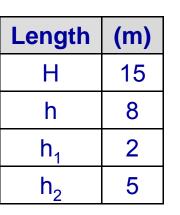
• For sloshing inside the tank of a LNG carrier or for the motion of a wave energy converter, numerical simulations can provide impressive results but the question remains of how relevant these results are when it comes to determining impact pressures !

• The numerical models are too simplified to reproduce the high variability of the measured pressures. NOT POSSIBLE FOR THE TIME BEING TO SIMULATE ACCURATELY BOTH GLOBAL AND LOCAL EFFECTS ! (see ISOPE 2009 Numerical Benchmark)

• However, numerical studies can be quite useful to perform sensitivity analyses in idealized problems (see ISOPE 2010 Numerical Benchmark)

COMPARATIVE NUMERICAL STUDY (2010)





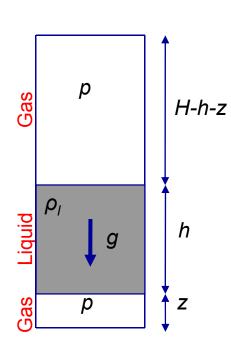
Organizers: GTT and UCD (*compressible bi-fluid software was required*) **1D case**

LNG = Liquefied Natural Gas NG = Natural Gas

Participants
ANSYS
Principia
ENS-Cachan
Hydrocean
Bureau Veritas
Lloyd's Register
Force

Case #	Scale	Liquid	Gas
1	1:1	LNG	NG
2	1:40	LNG	NG
3	1:40	Water	Air
4	1:40	Water	$SF_6 + N_2$
5	1:40	1:40-scaled LNG	1:40-scaled NG

1D SURROGATE MODEL OF AIR-POCKET IMPACT



Piston model

Perfect gas

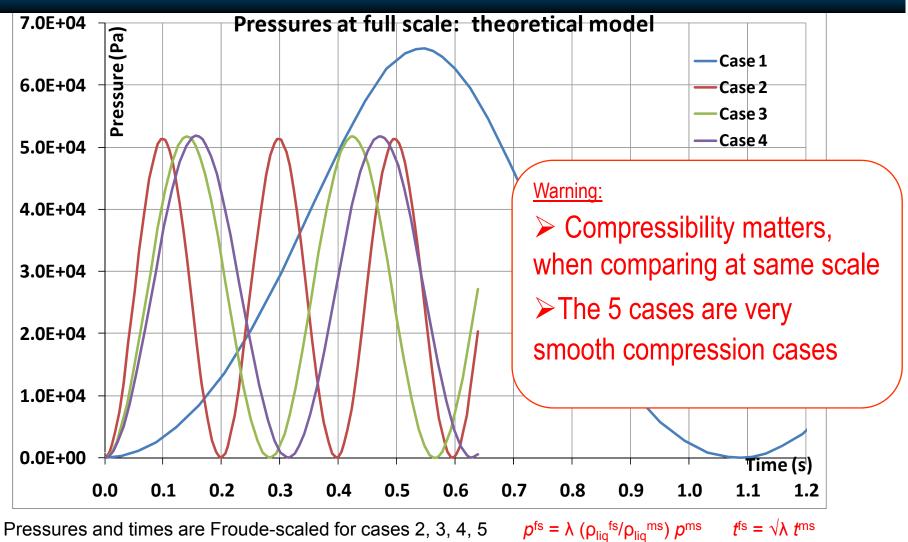
Adiabatic process

$$p = \rho RT$$
$$p \left(\frac{1}{\rho}\right)^{\gamma} = constant$$

Initial conditions: $p=p_0$,

$$z(0)=h_{l}, \quad \dot{z}(0)=0$$

THEORY : PRESSURES IN TIME DOMAIN



with λ = 40, ^{fs} = full scale, ^{ms} = model scale

CONCLUSIONS FOR THE 1D CASE

□ The different numerical methods are able to simulate adequately a simple smooth compression of a gas pocket without escape of gas

□ Very good agreement on the maximum pressure

□ For all methods :

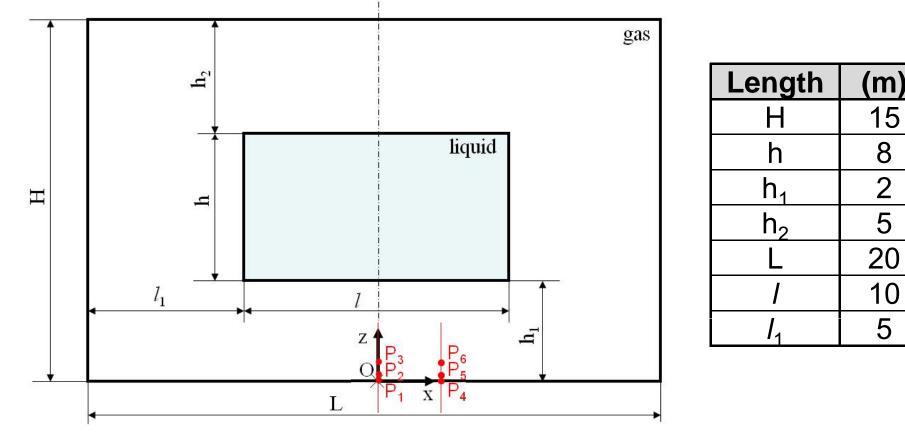
Complete Froude Scaling (CFS) works (same result for cases 1 & 5)

Partial Froude Scaling (PFS) generates a bias

COMPARATIVE NUMERICAL STUDY (2010)

VIDEO 2 – ENS-Cachan code





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COMPARATIVE NUMERICAL STUDY (2010) – 2D CASE

Case #	Scale	Liquid	Gas
6	1:1	LNG	NG
7	1:40	LNG	NG
8	1:40	Water	Air
9	1:40	Water	SF ₆ +N ₂
10	1:40	1:40-scaled LNG	1:40-scaled NG

Participants	Software	Method	6	7	8	9	10
ANSYS	Fluent	Finite Volume/VOF	Х	Х	Х	Х	Х
Principia	LS-DYNA	FEM Euler/Lagrange	Х	Х	Х	Х	Х
ENS-Cachan	Flux-IC	Finite Volume/NIP	Х	Х	Х	Х	Х
Lloyds Register	OpenFOAM	Finite Volume/VOF	Х	Х	Х	Х	Х
Force	Comflow	Finite Volume/VOF	Х	Х	Х	Х	Х
UoSFSI	in House	Finite Differences/VOF	Х	Х		Х	

THE ENS-CACHAN CODE

European Journal of Mechanics B/Fluids 28 (2009) 475-485



A totally Eulerian finite volume solver for multi-material fluid flows

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^c CMLA, ENS Cachan and CNRS, UniverSud and LRC MESO, ENS Cachan and CEA/DIF, 61 avenue du Président Wilson, F-94235 Cachan Cedex

ARTICLE INFO

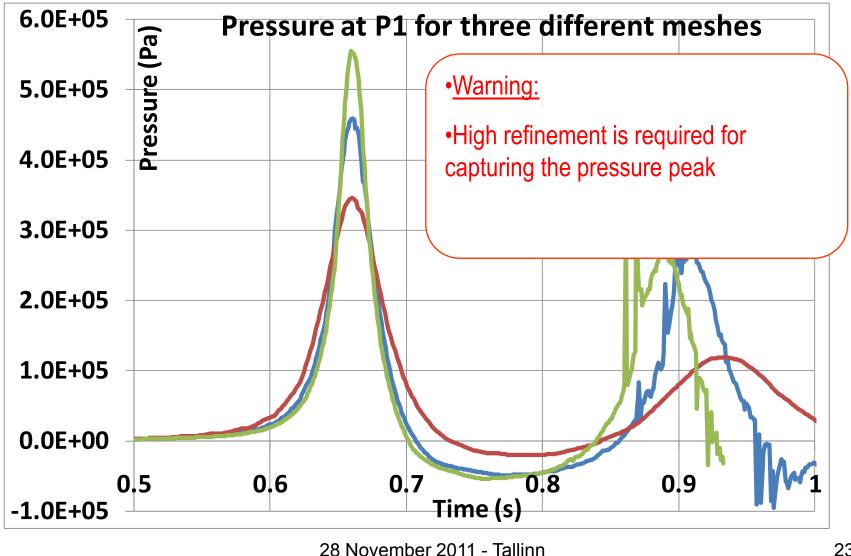
Article history: Received 28 April 2008 Received in revised form 19 March 2009 Accepted 19 March 2009 Available online 26 March 2009

Keywords: Compressible hydrodynamics Finite volume method

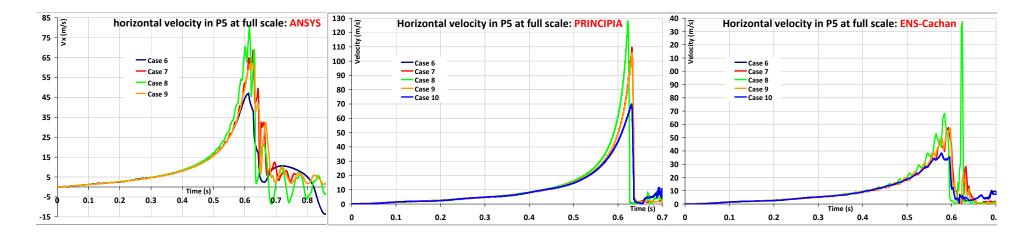
ABSTRACT

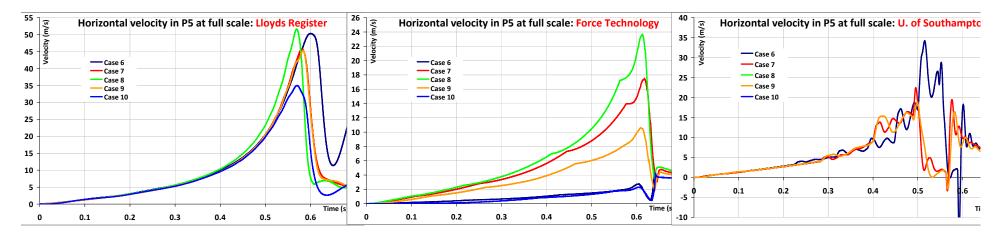
The purpose of this work is to present a new numerical scheme for multi-material fluid flows in dimension $d \ge 1$. It is a totally Eulerian conservative scheme that allows to compute sharp interfaces between non-miscible fluids. The underlying flux scheme in single material cells is the so-called FVCF scheme, whereas interface reconstruction and directional splitting is used in multi-material cells. One of the novelty of our approach is the introduction of the concept of "condensate" which allows to handle mixed cells containing two or more materials. Moreover, it has been designed to allow free sliding of materials on each others, thanks to a material volume centered computation of variables in mixed cells. @ 2009 Elsevier Masson SAS. All rights reserved.

ENS-CACHAN : MESH SENSITIVITY FOR CASE 6 (SCALE 1:1)



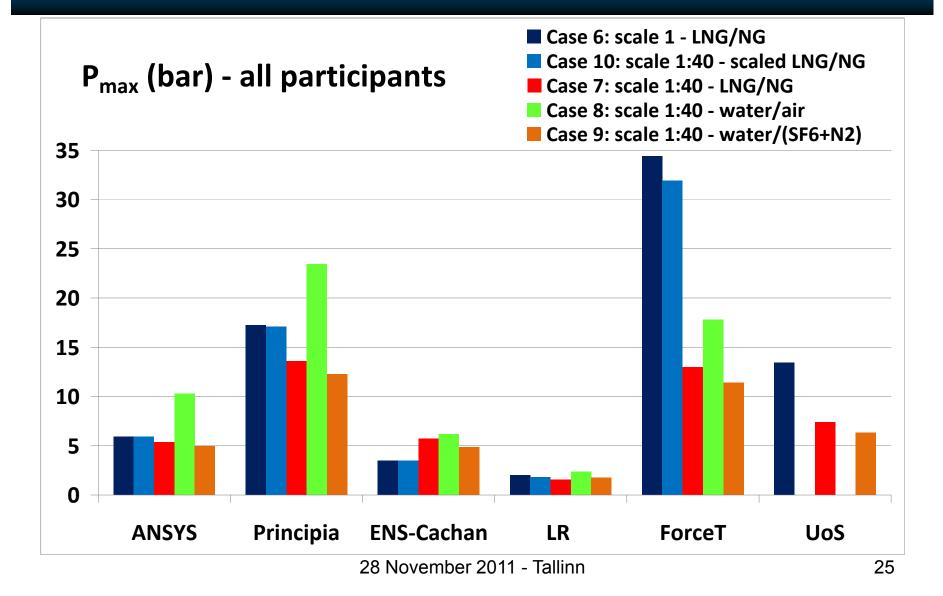
HORIZONTAL TIME SERIES AT P₅ : ALL PARTICIPANTS





Velocities and times are Froude-scaled for cases 7, 8, 9, 10 $V^{fs} = \sqrt{\lambda}$. V^{ms} , $t^{fs} = \sqrt{\lambda}$. t^{ms} with $\lambda = 40$, $f^{fs} = full scale$, $m^{ms} = model scale$ 28 November 2011 - Tallinn

MAXIMUM PRESSURE AT P₁ : ALL PARTICIPANTS



COMPARATIVE NUMERICAL STUDY (2010) – 2D CASE

 $\hfill\square$ Absolute values for V_{max} and P_{max} are very scattered

The meshes are not refined enough to capture sharp peak pressures
 After some work on the models, results should be much less scattered (work in progress)

□ Such a simple test should be passed adequately before

attempting to calculate more complex impacts

□ For all methods, whether relevant or not :

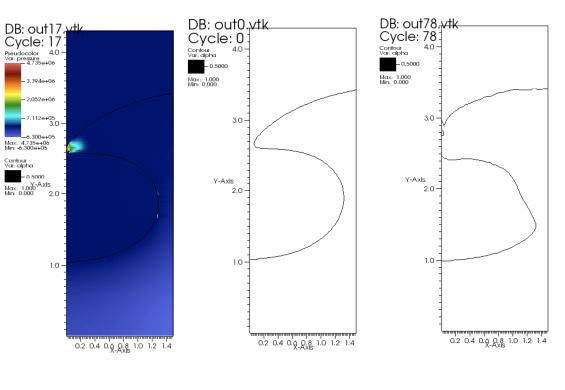
- Complete Froude Scaling (CFS) is satisfied
- Partial Froude Scaling generates a bias

BACK TO WAVE IMPACT

• Strategy used to compute wave impact : couple potential flow solver with two-fluid compressible flow solver

 Potential flow solver computes the wave all the way to overturning (Fochesato & Dias 2006)

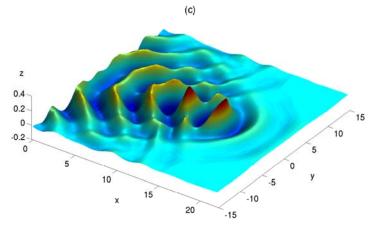
• Two-fluid (gas + liquid) compressible flow solver computes the liquid impact on the wall



work in progress

THE FULLY NONLINEAR POTENTIAL FLOW SOLVER

VIDEO 5 – Experiments in Nantes



PROCEEDINGS

Proc. R. Soc. A (2006) 462, 2715–2735 doi:10.1098/rspa.2006.1706 Published online 5 April 2006

A fast method for nonlinear three-dimensional free-surface waves

By Christophe Fochesato and Frédéric Dias*

□ Evidence of directional wave focusing in a « numerical » wave tank (Fochesato, Grilli, Dias, Wave Motion, 2007)

□ High-order three-dimensional boundary element method combined with mixed Eulerian–Lagrangian time updating, based on second-order explicit Taylor expansions with adaptive time-steps

□ Accelerated by the Fast Multipole Algorithm

OCEAN WAVE ENERGY : AN ASSET

	WAVE ENERGY CONVERSION
	Oyster Aquamarine Power 2009
	Oyster Aquamarine Power 2011

WAVE ENERGY CONVERSION

Aquamarine Power is a technology company that has developed a product called Oyster which produces electricity from ocean wave energy.

> UCD and Aquamarine Power are collaborating to deliver the nextgeneration Oyster 800.

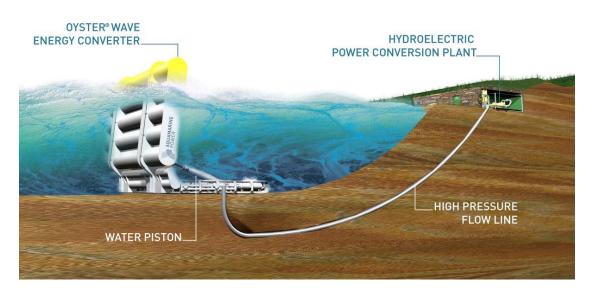




OYSTER TECHNOLOGY

- Large mechanical 'flap' moves back and forth with motion of waves
- Two hydraulic pistons pump high pressure water via pipeline to shore
- Conventional hydroelectric generator located onshore
- Secured to seabed at depths of 8 – 16m
- Located near shore, typically 500
 – 800m from shoreline

VIDEO 7





- Oyster 1 Project 315kW demonstrator successfully installed and grid-connected at European Marine Energy Centre (EMEC) in Orkney, October 2009 – Spring 2011 (finished)
- Oyster 2 Project 2.4MW project (3 Oyster 800 WEC) on schedule for 2011(1 Oyster 800 installed) / 2012 (two more to be installed)
- Oyster 3 Project 10MW development on track commissioning 2013 or 2014
- First commercial wave farm off the west coast of Ireland (Westwave project) 2015



WHAT DOES MATHEMATICS BRING ? High end computational modeling for wave energy systems

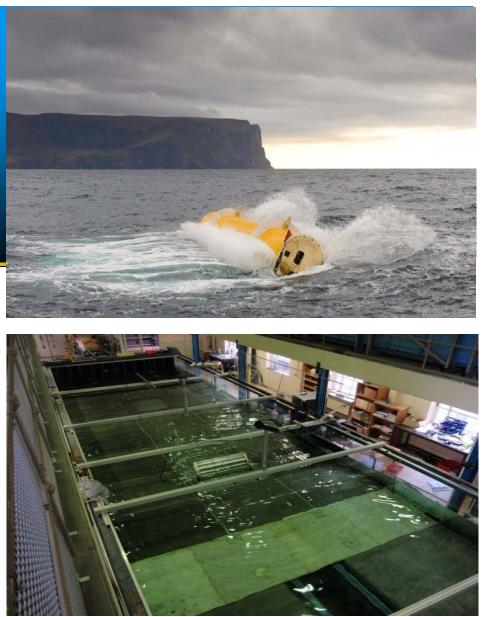


- 1. Wave impact and pressure loads on a single Wave Energy Converter
- 2. Optimal device spacing for an array of Wave Energy Converters
- 3. Preferred geographical locations for near shore Wave Energy Converter sites in Ireland
- 4. Biofouling (biological growth on surfaces in contact with water)



SURVIVABILITY Scale effects in experiments at small scale

- One of the most commonly acknowledged difficulties of conducting experiments with Wave Energy Converters : presence of scale effects (Reynolds much larger at full scale than at small scale – for example, at scale 1/40, viscous forces on the model are multiplied by a factor 253 if only Froude scaling is satisfied)
- This makes mathematical and numerical modelling a particularly valuable tool in the development of Wave Energy Converters



Experiments performed at Queen's University Belfast



WHICH TOOLS ? High end computational modeling for wave energy systems



Wave impact and pressure loads on a single Wave Energy Converter

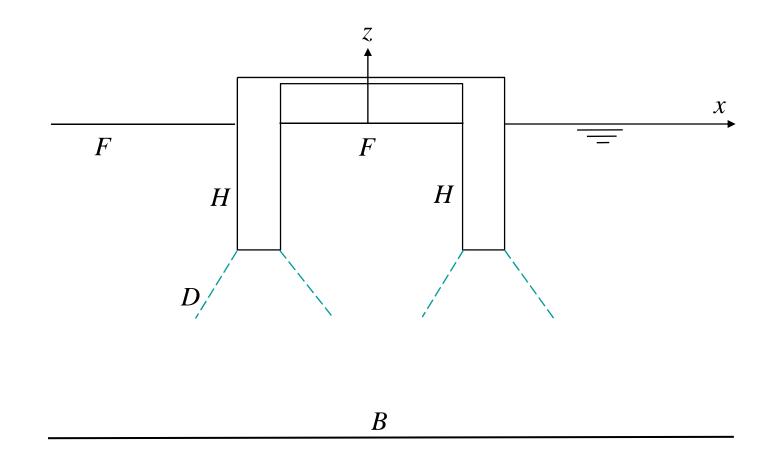
Fully nonlinear potential flow solver (for example WSIM – Boundary Element Method) combined with a Navier-Stokes solver (for example OpenFOAM – vortex shedding at the edges of the flap)

Alternative approach (dissipative surfaces)

Oyster 800 prototype is equipped with a number of pressure sensors

Fairly perfect fluid (joint work with X.B. Chen – Bureau Veritas)

Fluid domain with a body and dissipative surfaces





WHICH TOOLS ? High end computational modeling for wave energy systems

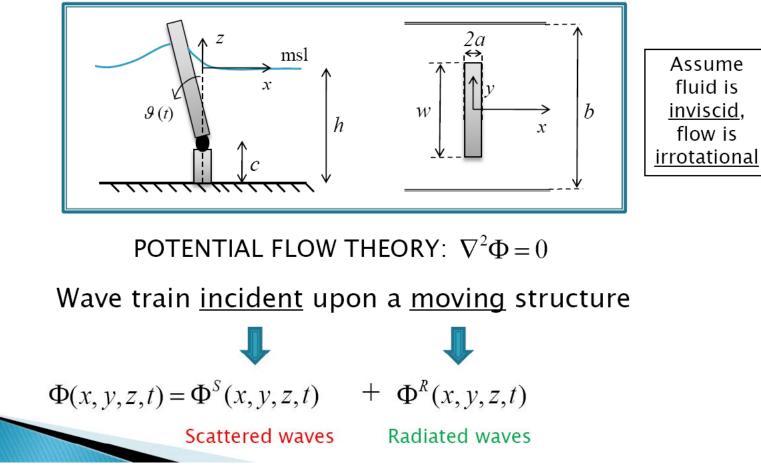


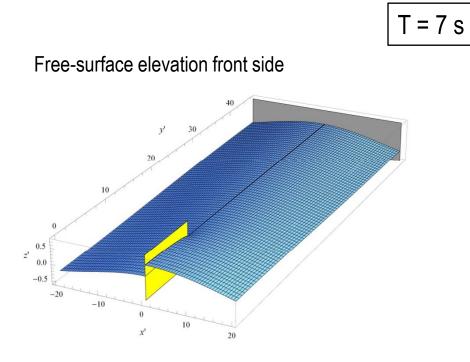
Optimal device spacing for an array of Wave Energy Converters

Analytical approach for the time being

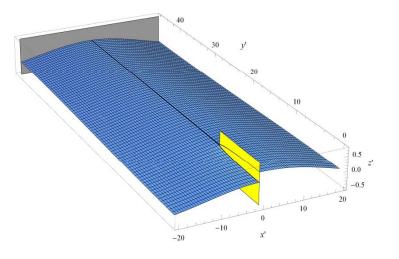
PRELIMINARY WORK ON ARRAYS OF FLAPS (E. Renzi)

3D LINEAR THEORY

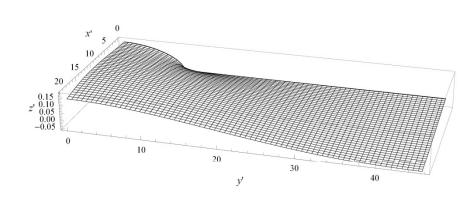




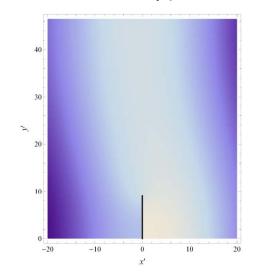
Free-surface elevation back side



Reflected waves



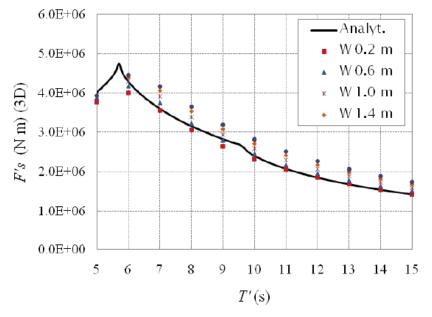
Free-surface density plot



PRELIMINARY WORK ON ARRAYS OF FLAPS

3D LINEAR THEORY

Scattering problem - Validation



OWSC width	18 m
Dist. btm-hinge	1.5 m
Water depth	10.9 m
Ampl. inc. wave	0.3 m

- Comparison with numerical model (WAMIT data by J. van't Hoff) very satisfactory (time mesh should be refined)
- Thin-plate approximation validated (good also for thicker plate)

Analytical theory for Oyster

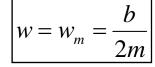


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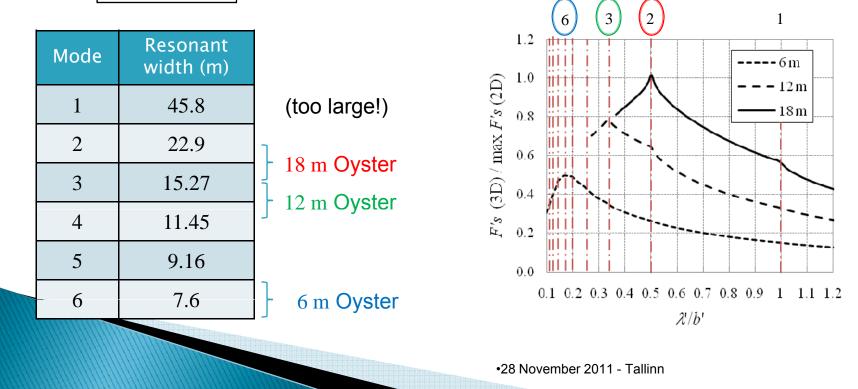
3D LINEAR THEORY

Scattering problem - Discussion

For a given modal order m, the modal interaction factor is max when



RESONANT WIDTH FOR MODE *m*



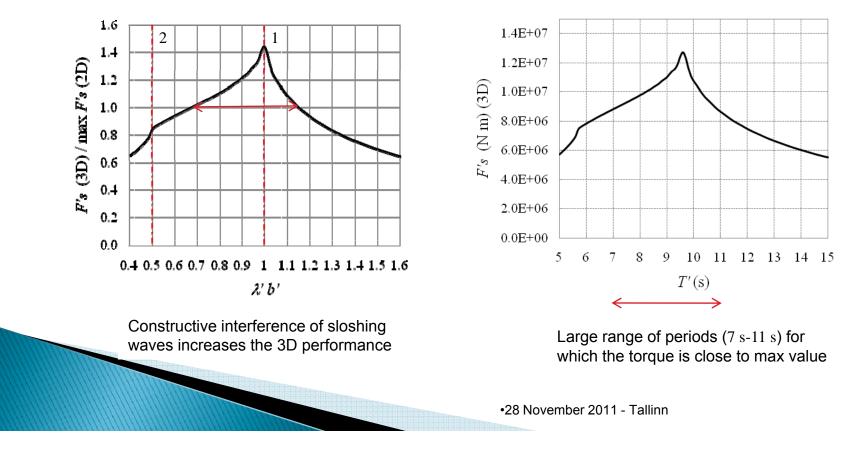
Analytical theory for Oyster



3D LINEAR THEORY

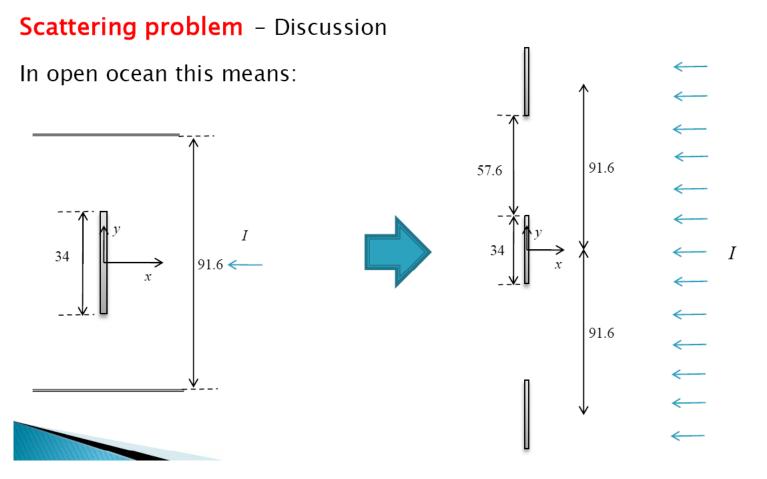
Scattering problem - Discussion

An optimum width (in <u>this</u> channel!) for which the Oyster interacts mostly with the first two sloshing modes: $w_{opt} \approx (46+23)/2 \text{ m} \approx 34 \text{ m}$



PRELIMINARY WORK ON ARRAYS OF FLAPS

3D LINEAR THEORY



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WHICH TOOLS ? High end computational modeling for wave energy systems



Preferred geographical locations for near shore Wave Energy Converter sites in Ireland

Variety of spectral wave models available to scientists and engineers, including SWAN and WAVEWATCH III[™] (one solves the random phase spectral action density balance equation for wavenumber-direction spectra). Is it enough for wave prediction or is it necessary to couple such spectral wave models with shallow-water type models in very shallow water?

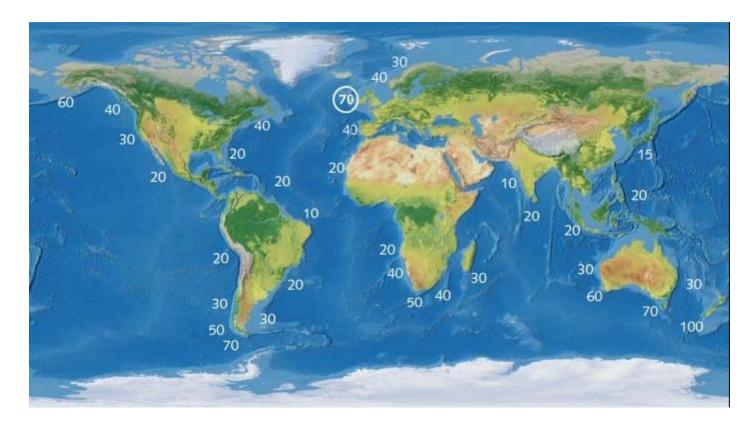


WHICH TOOLS ? High end computational modeling for wave energy systems



Biofouling (biological growth on surfaces in contact with water)

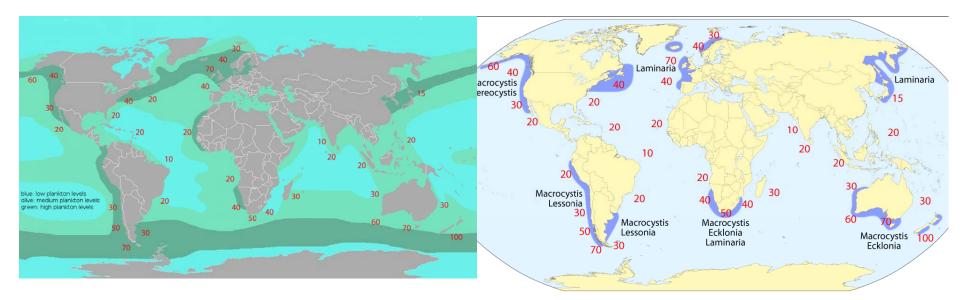
WAVE POTENTIAL AROUND THE WORLD



Numbers show the available power in kW/m (Source : World Energy Council)

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WAVE POTENTIAL + PLANKTON AND KELP DISTRIBUTIONS



Source : Kelp data : Kelp forest ecosystems - biodiversity, stability, resilience and future, Robert S. Steneck, Michael H. Graham et al.

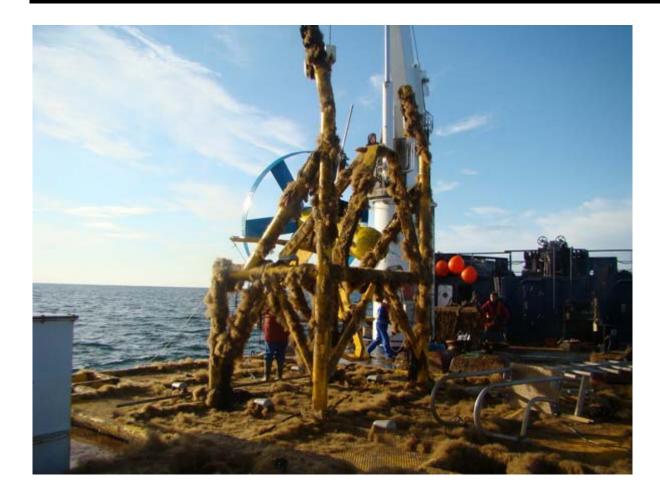
Temperate Shelf and Seas

seasonal variability + freshwater influx from coastal streams and tidal action

- + very heterogeneous habitats
- = high diversity of organisms: algae, invertebrate, fish, marine mammals,...

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BIOFOULING MUST BE CONSIDERED SERIOUSLY

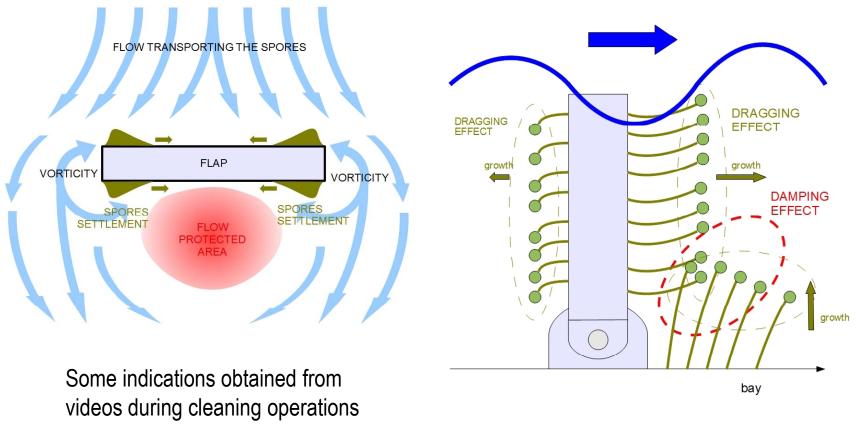


A tidal device which spent one year in water on Ouessant island in Britanny ...

> Modeling of growth done in collaboration with E. Reynaud (biologist at UCD) and C. Pinck

work in progress

EXPECTED SETTLEMENT ON THE FLAP



Mechanical effects on the flap

CONCLUDING REMARKS

Flaps might be simple devices from the engineering point of view. But they are really **challenging** from the fluid mechanics and numerical simulation points of view

- □ Intermediate water depth (kh of order unity)
- □ 3D problem
- □ Neither laminar flow nor potential flow (vortex shedding)
- □ Moving solid boundary (large motion) + free surface
- □ Intermediate size structure (no simplification)
- □ Nonlinear waves
- Coupling with biological growth
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THANK YOU FOR YOUR ATTENTION



Howth, Ireland